Charge transport model for photovoltaic devices based on printed polymer: Fullerene nanoparticles

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A R T I C L E   I N F O

Article history:
Received 19 March 2015
Received in revised form
22 May 2015
Accepted 24 May 2015
Available online 11 June 2015

Keywords:
Electrical transport model
Polymer nanoparticle
Aqueous dispersion
Flexible solar cells
Doctor blade coating

A B S T R A C T

The electrical transport properties of films derived from aqueous semiconducting nanoparticles are fully described by a phenomenological model that relates intrinsic film morphology to photovoltaic response. The model is applied to a new bulk heterojunction blend, composed of organic semiconducting nanoparticles formed from the polymer donor poly[(2,6-(4,8-didodecylbenzo[1,2-b:4,5-b']dithiophene))-alt-[5,5-(2,5-bis(2-butyloctyl)-3,6-dithiophen-2-yl-2,5-dihydropyrrolo[3,4-c][57]pyrrole-1,4-dione)] (P(TBT-DPP)) and the fullerene indene-C60-bisadduct (ICBA) synthesized by the miniemulsion method. The nanoparticle inks are printed from an aqueous dispersion onto flexible ITO-free substrates yielding power conversion efficiency of 2.6%.

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1. Introduction

Solar cells based on organic semiconductors have attracted significant interest due to their potential for large area production, lightweight and new form factors [1,2]. These devices are often fabricated using print[3–8] and coating techniques that are compatible with flexible substrates enabling progress towards low-cost high speed production methods [9–11]. The optimization of bulk heterojunction solar cells based on polymer:fullerene blends lead to recent reports of high performance, reaching up to 10% power conversion efficiencies (PCE) [12–14]. The high PCEs achieved with devices based on blends of electron-donating and electron-accepting materials require fine control of thin film phase separation on the nanometer scale which leads to increased number of interfaces where exciton dissociation and charge generation occurs [15]. The resulting film morphology strongly depends on several processing parameters [16] including the solubility of the materials in the solvent used [17], the interaction with the substrate surface [18] the packing of the molecules, formation of domains of different blend compositions [19] and the film deposition technique [20]. It has been established that exciton dissociation can be maximized when donor and acceptor components are intimately mixed and the length-scale of the phase separation is in the range of the exciton diffusion length [15,21]. As a result, over the years, many methodologies to control the blend morphology have been developed: appropriate solvent choice [22] controlling the rate of solvent evaporation [23,24], thermal annealing [25,26] melting of bilayers [27] and adding solvent additives [28]. These different procedures promote the formation of a phase separated morphology that leads to improved photovoltaic performance. However the prediction and control of the final film morphology is extremely challenging due to the interplay between mixing and drying. Most laboratories use toxic and flammable organic solvents in their fabrication process of organic solar cells. Polymer nanoparticles dispersed in water are an alternative method that addresses both morphology control and scalable on non-toxic ink deposition. The miniemulsion technique [29] retains the donor/acceptor phase segregation achieved at the synthesis step regardless of the deposition solvent evaporation time or film processing method. This technique was previously utilized for synthesis of blends of polyfluorenes [30,31] and poly(3-hexylthiophene)(P3HT); phenyl-C60-butyric acid methyl ester (PCBM) nanoparticles [32] deposited by blade coating achieving PCEs of 0.29% [33] and 0.55% [34]. Solar cells based on spin coated blend nanoparticles of P3HT and the fullerene indene-C60-bisadduct (ICBA) resulted in power conversion efficiency of 2.50% [35]. Although several photovoltaic devices based on aqueous nanoparticle dispersions have been reported [30–32,35,36], the effects of
Fig. 1. (a) Simplified scheme of the nanoparticle synthesis: organic solvent solution containing P(TBT-DPP) and ICBA is added to an aqueous solution containing surfactant, this mixture is sonicated to form the nanoparticle dispersion and after a dialysis process the nanoparticle ink is ready for printing; (b) scheme describing the doctor blade coating technique; (c) device structure: PEN/PEDOT:PSS/NP P(TBT-DPP):ICBA/C60/Al; (d) device photograph and (e) equivalent circuit used to model the dark J–V response of the devices: \( R \) is the series resistance, \( R_{sh} \) is the shunt resistance, \( V_c \) is the potential drop across \( R \), \( J_D \) is the diode current, \( J_{sh} \) is the shunt current and \( f \) is the total current of the system.

this type of donor/acceptor phase segregation on the electrical properties remains poorly explored. We fabricated blade coated devices using conventional chloroform bulk heterojunction (BHJ) and nanoparticle dispersed in aqueous media (NP) to form the active layer of flexible solar cells. We characterized the film morphology and electrical properties and fitted the experimental dark \( J–V \) curves using a phenomenological electric model. The nanoparticle composition and morphology were correlated to the photovoltaic performance and we demonstrate that the charge transport in nanoparticulate devices is improved when compared to the BHJ. The photovoltaic active layer of the flexible solar cells reported here, is composed of a new nanoparticulate bulk heterojunction blend of the polymer poly[2,6-(4,8-diododecylbenzo[1,2-b:4,5-b’]dithiophene)-alt-(5,5-(2,5-bis(2-butylyloxy)-3,6-dithiophen-2-yl-2,5dihydropyrrolo[3][5]pyrrole-1,4-dione)] (P(TBT-DPP)) [37] with ICBA. The P(TBT-DPP):ICBA nanoparticle ink was synthesized in aqueous media following the miniemulsion technique illustrated in Fig. 1a. The aqueous nanoparticle (NP) inks were printed by doctor blade coating in air [38] onto PEN/PEDOT:PSS substrates as depicted in Fig. 1b and a schematic representation of the device structure is shown in Fig. 1c. Conventional bulk heterojunction (BHJ) devices were also deposited by blade coating and used as control devices. Our model based on [39,40] can be applied to fit the dark \( J–V \) curves of both conventional BHJ and nanoparticle devices. Our model assumes that the dark \( J–V \) curves can be obtained using the equivalent circuit illustrated in Fig. 1e. The model is composed of a nonlinear space-charge limited series resistance (\( R \)), that depends on \( J \), connected to a diode (\( D \)) in parallel to a shunt resistance (\( R_{sh} \)) [39,40]. By means of this Space-Charge Limited-Diode (SCL-D) model, relevant parameters such as effective mobility (\( \mu_{eff} \)) and the diode saturation current (\( J_0 \)) are quantified and related to the photovoltaic response of the devices.

2. Experimental section

Synthesis and characterization of the nanoparticle inks: dispersions were prepared using P(TBT-DPP) provided by Polyaera and ICBA (purchased from Sigma Aldrich). Chloroform solutions of P (TBT-DPP) and ICBA (40 mg/mL) in three donor:acceptor weight ratios (DAWR) (1:1, 3:7 and 1:3) were added to a 5 mg/mL aqueous sodium dodecyl sulfate (SDS) solution. Preemulsification is achieved by stirring the resulting solution at 1000 rpm during 1 h. Subsequently the mixture was sonicated using a Branson 450 digital sonifier for 2 min at 60% amplitude with a microtip (radius 6.5 mm) immersed in the emulsion to form the dispersion. Then the samples were heated while stirring for 30 min at 65 °C to evaporate the chloroform forming a stable aqueous dispersion. In order to concentrate and eliminate the excess surfactant, dialysis of the nanoparticle dispersion was performed using ultra centrifuge dialysis tubes with 10 kDa molecular weight cutoff (purchased from Millipore). The dialysis procedure was repeated until the surface tension of the filtrated water reached about 38 mN/m. Dynamic light scattering (Zetasizer Nano-ZS, Malvern Instruments) was used to measure the size distribution of the nanoparticles in the aqueous dispersion. The nanoparticle composition and average particle size measured by Dynamic Light Scattering are summarized in Table 1. The nanoparticle size strongly depends on the initial ratio of surfactant, polymer, fullerenic, organic solvent and water. Therefore the synthesis parameters can be tuned in order to control the final particle size. For a given blend ratio, higher amount of blend material leads to larger particle size [41] whereas the increase in the surfactant amount reduces the particle diameter [29]. In this work, the total blend concentration was kept constant at 40 mg/mL and it is shown that the average diameter increases with the amount of ICBA in the blend. The molecular weight of the polymer is 27 kDa [37]. Since the monomer molecular weight of the P(TBT-DPP) is approximately 1.12 kDa, the weight-average degree of polymerization (\( n_{wp} \)) is \( \sim 24 \). The molecular weight of ICBA, the acceptor, is 952.96 Da [42]. Based on the molecular weight values for donor and acceptor materials we can calculate the ratio (\( \eta \)) of ICBA molecules per polymer chain for each nanoparticle blend composition. The average density of ICBA molecules per monomer of the polymer inside the nanoparticle (\( \rho \)) is then given by \( \rho = \eta / V \), where \( V \) is the average volume of the nanoparticle (assuming a spherical structure with the diameter
3. Results and discussion

From Fig. 1(e), the current of the system is given by 
\[ J = J_{sh} + J_0 \exp\left(q(V - V_c)/kT\right) - 1 \]  
(1)

The system theoretical dark J–V characteristics can be obtained by solving Eq. (1) for a given \( J \), where \( J_0 \) and \( R_{sh} \) are fitting parameters. From the Mott–Gurney law, \( R_{sh} \) this is given by
\[ R_{sh} = \frac{8L^3}{9\mu_{eff}^2} \frac{e}{A} \frac{1}{V - V_c} \]  
(2)

where \( L \) is the sample thickness and \( \mu_{eff} \) is the fitting parameter that represents the effective mobility at low voltage. Applying the SCL-D model to the dark experimental J–V curves gives the charge carriers effective mobility (\( \mu_{eff} \)) which is also a fitting parameter that determines the space-charge limited (SCL) current at high voltages. The shunt resistance, \( R_{sh} \), controls the ohmic behavior at low voltages and \( J_0 \) is given by the steep increase in the current between the ohmic and the SCL regimes [39]. Both, \( J_0 \) and \( \mu_{eff} \) depend on the intrinsic morphology of the photoactive layer and can be directly related to the photovoltaic response of the device. For instance, small \( J_0 \) is associated with higher open circuit voltages by the expression \( V_{oc} = (kT/\mu_{eff}^2) \ln(J_{0}/J_0) \) [44]. Moreover the diode current can be written as [40]
\[ J_0 = J_{mat} \exp(-E_{DA}/2kT) \]  
(3)

where \( E_{DA} \) is the effective energy barrier at the donor/acceptor interface and \( J_{mat} \) is an independent factor determining the carrier recombination rate [40]. \( J_{mat} \) quantifies the inter-molecular overlap of the donor and acceptor molecules and the effective area of the donor/acceptor interface. By the same token, \( \mu_{eff} \) is also strongly affected by the morphology of the film at the molecular level. Large domains of donor (acceptor) aggregates create large barriers for hole (electron) transport. The increased volume of those phases tends then to decrease the number of percolation paths that link the dissociation centers to the metallic contacts. As a consequence, large domains of polymer (fullerene) would tend to block the electron (hole) transport toward the collecting electrodes. This effect increases the chances of recombination during free carrier drift that follows exciton dissociation. A higher rate of recombination decreases \( \mu_{eff} \) which results in lower short circuit currents. The film morphologies for compositions of 3:7 (DAWR) for chloroform BHJ and aqueous NP devices were accessed using atomic force microscopy (AFM) and the images are displayed in Fig. 2a and b, respectively and the phase contrast image simultaneously acquired are shown in Fig. 2c and d. The conventional BHJ film blade coated from chloroform solutions resulted in large domains in the order of 500 nm of diameter. In contrast, the P (TBT-DPP):ICBA nanoparticle film exhibited finer phase segregation on the order of 100 nm, which is consistent with the DLS measurement for the aqueous dispersion. These images confirm that blend phase segregation fixed during the nanoparticle synthesis is maintained during film deposition and solvent drying. Fig. 2e and f shows that there is a good agreement between the model proposed here and the experimental data obtained for both types of morphology. In Table 2 we report the values of \( \mu_{eff} \) and \( R_{sh} \) obtained from the model at V = 0.1 V. As displayed in Table 2, \( \mu_{eff} \) and \( J_0 \) increase for the devices made with higher ICBA density nanoparticle solutions. The larger domains of donor (acceptor) shown in the BHJ film produce higher fluctuations in the local density of ICBA molecules. Those fluctuations mask the dependence of \( \mu_{eff} \) and \( J_0 \) on the bulk density of ICBA that is only apparent in the NP devices where the mixing of the acceptor and donor materials is finer and more evenly distributed throughout the film. The values in Table 2 clearly show that the NP devices always give higher \( \mu_{eff} \) and \( J_0 \) when compared to the respective BHJ cells for the same DAWR.
The $J-V$ curves under illumination for conventional BHJ and NP devices in the blend DAWR of 1:1, 3:7 and 1:3 are shown in Fig. 3 and all the photovoltaic parameters obtained from these curves are summarized in Table 3.

It was observed in similar polymer/acceptor nanoparticulate photovoltaic devices that NP materials show evidence of a higher inter-chain order compared to the BHJ films [32]. This order tends to increase the effective mobility for the NP devices, increasing $J_{\text{mat}}$ given in Eq. (3) which enhances $J_0$. In contrast, the lower level of control over the nanoscale morphology in BHJ devices would reduce the effective donor/acceptor heterojunction area, leading to a lower $J_{\text{mat}}$. As a consequence, the NP devices have higher $J_0$ when compared to the BHJ devices for the same DAWR. Higher $J_0$ translates into lower values of $V_{\text{oc}}$ [44]. The NP devices generate lower open circuit voltages under illumination (100 mW cm$^{-2}$) when compared to the respective BHJ devices as it is shown in Fig. 2.

Table 2

<table>
<thead>
<tr>
<th>Blend weight ratio</th>
<th>$J_0$ [mA cm$^{-2}$]</th>
<th>$\mu_{\text{eff}}$ [cm$^2$ V$^{-1}$ s$^{-1}$]</th>
<th>$R_{\text{sh}}$ [Ω cm$^2$]</th>
<th>$V_{\text{oc}}^{\text{exp}}$ [V]</th>
<th>$V_{\text{oc}}^{\text{calc}}$ [V]</th>
</tr>
</thead>
<tbody>
<tr>
<td>BHJ (1:1)</td>
<td>$3 \times 10^{-8}$</td>
<td>$6.5 \times 10^{-14}$</td>
<td>$1.43 \times 10^2$</td>
<td>$0.52 \pm 0.09$</td>
<td>$0.45$</td>
</tr>
<tr>
<td>BHJ (3:7)</td>
<td>$1 \times 10^{-9}$</td>
<td>$5.0 \times 10^{-10}$</td>
<td>$5.66 \times 10^3$</td>
<td>$0.55 \pm 0.01$</td>
<td>$0.56$</td>
</tr>
<tr>
<td>BHJ (1:3)</td>
<td>$2 \times 10^{-7}$</td>
<td>$1.0 \times 10^{-11}$</td>
<td>$3.34 \times 10^3$</td>
<td>$0.42 \pm 0.01$</td>
<td>$0.42$</td>
</tr>
<tr>
<td>NP (1:1)</td>
<td>$4 \times 10^{-8}$</td>
<td>$7.0 \times 10^{-12}$</td>
<td>$7.90 \times 10^2$</td>
<td>$0.45 \pm 0.02$</td>
<td>$0.47$</td>
</tr>
<tr>
<td>NP (3:7)</td>
<td>$1 \times 10^{-9}$</td>
<td>$1.0 \times 10^{-12}$</td>
<td>$1.66 \times 10^3$</td>
<td>$0.57 \pm 0.01$</td>
<td>$0.55$</td>
</tr>
<tr>
<td>NP (1:3)</td>
<td>$2 \times 10^{-8}$</td>
<td>$8.0 \times 10^{-12}$</td>
<td>$1.43 \times 10^4$</td>
<td>$0.49 \pm 0.03$</td>
<td>$0.48$</td>
</tr>
</tbody>
</table>

Table 2. This behavior was also reported when other NP systems were compared to the corresponding BHJ [32,36]. The higher density $\rho$ of ICBA molecules per monomer of polymer in the NP
A core–shell structure formed of an acceptor rich core surrounded by a polymer-rich shell was previously observed in other nanoparticulate donor/acceptor films [32,37]. Since in those structures the acceptor is reasonably dispersed inside the polymeric shell, the transport would be essentially limited by the hopping of electrons across thin polymeric layers coating the adjacent ICBA molecules. Assuming that the polymer acts as a potential barrier to inter-acceptor electron hopping, the transport of electrons between adjacent electronic states localized in the acceptor at constant temperature may depend only on the tunneling part of the hopping rate. It is then reasonable to assume that $\mu_{\text{eff}} \propto -d$ [45], where $d$ is the average distance between neighbor ICBA molecules in the NP. Due to spatial considerations and neglecting the volume of the fullere molecule compared to the volume of the NP’s, one can write $\mu_{\text{eff}}(4\pi/3)d/2 \propto V$. It follows then that $d \propto \rho^{-1/3}$. This suggests that the relation $\mu_{\text{eff}} \propto -\rho^{-1/3}$ is valid for the NP structures used in our study. We use this relation and plot in Fig. 4 the effective mobilities from Table 2 as a function of the respective values of $\rho$ taken from Table 1. The predicted linear behavior is observed, with the solid line as a fit to the $\mu_{\text{eff}}$ data. This result indicates that the transport in the NP samples are limited by the inter-acceptor hopping of electrons across the barriers formed by the polymeric coating due to the finer donor/acceptor mix in the NP films.

Following the analysis above, the reasons behind the best photovoltaic performance obtained for the 3:7 NP solar cell are better understood. First the substantial decrease of the series resistance exponentially increases the effective mobility of the 3:7 NP solar cell. Additionally, the higher concentration of donor/acceptor heterojunctions in this sample (higher $\rho$) increases the density of photoexcited excitons upon illumination which enhances $J_{\text{sc}}$. Both phenomena compensate the lower $V_{\text{oc}}$ resulting in the highest PCE compared to the conventional BHJ devices we have analyzed.

Given the good performance shown by the devices with DAWR of 3:7, we further optimized the ink formulation to improve wettability and promote better film formation. 20% of ethanol in volume was added to the initial 3:7 to decrease the surface tension [46] of the nanoparticle aqueous dispersions. The addition of ethanol to the nanoparticle ink reduces the contact angle between ink and substrate, PEN/PEDOT:PSS, from 59° to 45° leading to a more homogeneous film coating. Fig. 5 displays the $J$–$V$ curve under illumination of 100 mW cm$^{-2}$ for the optimized NP device and the conventional BHJ in the same DAWR and blade coated in the same conditions. The NP device exhibited six fold increase in the $J_{\text{sc}}$ (2.12 mA cm$^{-2}$ for BHJ and 12.73 mA cm$^{-2}$ for NP) enhancing the PCE from 0.44% to 2.63%. This result is comparable to the highest PCE reported for P3HT:ICBA nanoparticles dispersed in water and spin coated onto ITO/PEDOT:PSS substrates (2.50%).

### Table 3

<table>
<thead>
<tr>
<th>Blend weight ratio</th>
<th>Active layer thickness [nm]</th>
<th>$J_{\text{sc}}$ [mA cm$^{-2}$]</th>
<th>$V_{\text{oc}}$ [V]</th>
<th>FF [%]</th>
<th>$R_{\text{s}}$ [$\Omega$ cm$^{2}$]</th>
<th>PCE [%]</th>
</tr>
</thead>
<tbody>
<tr>
<td>BHJ (1:1)</td>
<td>140 ± 10</td>
<td>1.17 ± 0.21</td>
<td>0.52 ± 0.09</td>
<td>30 ± 6</td>
<td>285 ± 72</td>
<td>0.18 ± 0.03</td>
</tr>
<tr>
<td>NP (1:1)</td>
<td>130 ± 10</td>
<td>2.14 ± 0.24</td>
<td>0.45 ± 0.02</td>
<td>45 ± 6</td>
<td>47 ± 11</td>
<td>0.44 ± 0.05</td>
</tr>
<tr>
<td>BHJ (3:7)</td>
<td>120 ± 8</td>
<td>2.12 ± 0.08</td>
<td>0.55 ± 0.01</td>
<td>37 ± 3</td>
<td>83 ± 10</td>
<td>0.44 ± 0.02</td>
</tr>
<tr>
<td>NP (3:7)</td>
<td>110 ± 15</td>
<td>2.53 ± 0.13</td>
<td>0.42 ± 0.01</td>
<td>52 ± 5</td>
<td>31 ± 7</td>
<td>0.56 ± 0.05</td>
</tr>
<tr>
<td>BHJ (1:3)</td>
<td>150 ± 5</td>
<td>1.80 ± 0.27</td>
<td>0.57 ± 0.01</td>
<td>44 ± 3</td>
<td>71 ± 17</td>
<td>0.44 ± 0.06</td>
</tr>
<tr>
<td>NP (1:3)</td>
<td>140 ± 12</td>
<td>2.37 ± 0.18</td>
<td>0.49 ± 0.03</td>
<td>44 ± 5</td>
<td>64 ± 13</td>
<td>0.52 ± 0.08</td>
</tr>
<tr>
<td>NP[3:7]</td>
<td>100 ± 26</td>
<td>10.49 ± 2.65</td>
<td>0.43 ± 0.01</td>
<td>47 ± 5</td>
<td>12 ± 3</td>
<td>2.16 ± 0.50</td>
</tr>
<tr>
<td>NP[3:7]</td>
<td>100 ± 26</td>
<td>12.73</td>
<td>0.44</td>
<td>47</td>
<td>11</td>
<td>2.63</td>
</tr>
</tbody>
</table>

* Optimized NP devices with the addition of 20% of ethanol in the NP ink.

* Best performing device (NP 3:7 with 20% EtOH), values taken from Fig. 5.
As it is shown in Fig. 5, despite the decrease in the $V_{oc}$, 0.55–0.44 V, the NP device has shown remarkable improvement in the fill factor (FF): from 37% (BHJ) to 46% (NP). This result is consistent with the decrease in the series resistance ($R_s$): from 83 Ω cm$^2$ (BHJ) to 11 Ω cm$^2$ (NP). We attribute the improved photovoltaic response found for the NP devices to the nanostructured film morphology.

4. Conclusions

We have fully described the transport properties of P(TBT-DPP):ICBA nanoparticles devices applying the SCL-D model to fit the dark J–V curves and demonstrate that the charge transport is improved for the NP when compared to conventional BHJ devices due to the closer contact between donor/acceptor in the nanostructured film. The best blend ratio was found to be (3:7) with optimized devices yielding power conversion efficiency of 2.63%. The aqueous nanoparticle ink significantly reduces the use of hazardous solvents in organic electronics. In addition, it allows control of the donor/acceptor morphology at the nanometer scale without compromising reproducibility in large area devices, such as photovoltaics. Our results confirm the viability of applying aqueous inks based on polymer:fullerene nanoparticles for large-area photovoltaics that are compatible with ITO-free flexible substrates and printing techniques.

Acknowledgments

This work is based upon work supported in part by the National Science Foundation under Cooperative Agreement no. ECC-1160494. Any opinions, findings and conclusions or recommendations expressed in this material are those of the author(s) and do not necessarily reflect the views of the National Science Foundation. The authors would also like to thank Polyera for providing materials and CNPq and CAPES–PDSE (process BEX – 1564129) for partial financial support.

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